## Heat equation with Neumann boundary conditions

We seek to solve

$$\rho \ c \ \frac{\partial u}{\partial t} - \kappa \ \frac{\partial^2 u}{\partial x^2} = f(x,t)$$
$$\frac{\partial u}{\partial x}(0,t) = \frac{\partial u}{\partial x}(l,t) = 0$$
$$u(x,0) = \psi(x)$$

## A preliminary problem

To generate some insight into the special behavior of this problem, we start by considering a simpler, related BVP:

$$-\kappa \frac{\mathrm{d}^2 u(x)}{\mathrm{d}x^2} = f(x)$$
$$\frac{\mathrm{d}u}{\mathrm{d}x}(0) = \frac{\mathrm{d}u}{\mathrm{d}x}(l) = 0$$

To handle this we introduce a differential operator  $L_N$  defined on the subspace of functions that satisfy a Neumann boundary condition.

$$L_N(u)=\{ ext{ -}\kappa \,\, u^{\prime\prime}\mid u ext{ such that } u^\prime(0)=u^\prime(l)=0 \;\}$$

The eigenfunctions of this differential operator are  $u_0(x) = 1$  and  $u_k(x) = \cos(n \pi x/l)$  for k = 1, 2, 3, ...

As before, we start by multiplying both sides of the ODE by these eigenfunctions and integrating over the interval in question. There are two separate cases to consider. For k = 0 we have

$$\frac{1}{l}\int_0^l -\kappa \, \frac{\mathrm{d}^2 u}{\mathrm{d}x^2} \, dx = \frac{1}{l}\int_0^l f(x) \, dx$$

The integral on the left simplifies to

$$\frac{1}{l}\int_0^l -\kappa \, \frac{\mathrm{d}^2 u}{\mathrm{d}x^2} \, dx = -\frac{\kappa}{l} \, \left(\frac{\mathrm{d}u}{\mathrm{d}x}\right) \Big|_0^l = 0$$

In general, the integral on the right won't be 0. That means that we are potentially stuck, unless we introduce a *compatibility condition*:

$$\frac{1}{l}\int_0^l f(x) \, dx = 0$$

The reason that we have to impose this extra condition is that the operator  $L_N(u)$  has a non-trivial null space. That means that we should expect the equation

$$L_N(u) = f$$

to not have a solution for some functions f(x). The work-around for the moment is to impose this extra condition on f(x) to guarantee that the problem will have a solution.

Assuming that this extra condition is satisfied, we can go to handle the integrals arising from the other eigenfunctions:

$$\frac{2}{l}\int_{0}^{l} -\kappa \frac{\mathrm{d}^{2}u}{\mathrm{d}x^{2}} \cos\left[\left(\frac{k\pi}{l}x\right)\right] dx = \frac{2}{l}\int_{0}^{l} f(x) \cos\left[\left(\frac{k\pi}{l}x\right)\right] dx$$

On the left we get after a couple of integrations by parts:

$$\frac{2}{l}\int_{0}^{l} -\kappa \frac{\mathrm{d}^{2}u}{\mathrm{d}x^{2}} \cos\left(\frac{k\pi}{l}x\right) dx$$
$$= -2 \frac{\kappa}{l} \frac{\mathrm{d}u}{\mathrm{d}x} \cos\left(\frac{k\pi}{l}x\right) |_{0}^{l} - \frac{2\kappa}{l^{2}} \frac{\kappa\pi}{l^{2}} \int_{0}^{l} \frac{\mathrm{d}u}{\mathrm{d}x} \sin\left(\frac{k\pi}{l}x\right) dx$$
$$= 0 - \frac{2\kappa}{l^{2}} \frac{\kappa\pi}{l^{2}} \left(u(x) \sin\left(\frac{k\pi}{l}x\right)\right) |_{0}^{l} + \frac{2\kappa}{l^{3}} \frac{\kappa^{2}\pi^{2}}{l^{3}} \int_{0}^{l} u(x) \cos\left(\frac{k\pi}{l}x\right) dx$$

## Solving the full problem

We now return to

$$\rho \ c \ \frac{\partial u}{\partial t} - \kappa \ \frac{\partial^2 u}{\partial x^2} = f(x,t)$$
$$\frac{\partial u}{\partial x}(0,t) = \frac{\partial u}{\partial x}(l,t) = 0$$
$$u(x,0) = \psi(x)$$

Before we seek to solve this, we may want to impose a compatibility condition similar to what we imposed in the earlier problem.

$$\frac{1}{l}\int_0^l f(x,t) \, dx = 0$$

If the f(x,t) we are given does not satisfy this compatibility condition, we can try the following trick. Introduce the function

$$g(t) = \frac{1}{l} \int_0^l f(x,t) \, dx$$

and let  $\beta(t)$  be the solution to

$$ho \ c \ eta'(t) = g(t)$$
 $ho(0) = rac{1}{l} \int_0^l \psi(x) \ dx$ 

and instead try to solve

$$\rho \ c \ \frac{\partial u}{\partial t} - \kappa \ \frac{\partial^2 u}{\partial x^2} = f(x,t) - g(t)$$
$$\frac{\partial u}{\partial x}(0,t) = \frac{\partial u}{\partial x}(l,t) = 0$$
$$u(x,0) = \psi(x) - \beta(0)$$

By construction, this modified problem will automatically satisfy the compatibility condition:

$$\frac{1}{l} \int_0^l f(x,t) - g(t) \, dx = \frac{1}{l} \int_0^l f(x,t) \, dx - \frac{1}{l} \int_0^l g(t) \, dx$$
$$= g(t) - \frac{1}{l} (l \ g(t)) = 0$$

How does the solution to the modified problem help us to solve the original problem? Let  $u_g(x)$  be the solution to

$$\rho \ c \ \frac{\partial u_g}{\partial t} - \kappa \ \frac{\partial^2 u_g}{\partial x^2} = f(x,t) - g(t)$$
$$\frac{\partial u_g}{\partial x}(0,t) = \frac{\partial u_g}{\partial x}(l,t) = 0$$
$$u_g(x,0) = \psi(x) - \beta(0)$$

by linearity of the operator we have that

$$\rho \ c \ \frac{\partial(u_g(x,t) + \beta(t))}{\partial t} - \kappa \ \frac{\partial^2(u_g(x,t) + \beta(t))}{\partial x^2}$$
$$= \left(\rho \ c \ \frac{\partial u_g}{\partial t} - \kappa \ \frac{\partial^2 u_g}{\partial x^2}\right) + \left(\rho \ c \ \frac{\partial(\beta(t))}{\partial t} - \kappa \ \frac{\partial^2(\beta(t))}{\partial x^2}\right)$$
$$= (f(x,t) - g(t)) + g(t) = f(x,t)$$
$$(u_g(x,0) + \beta(0)) = \psi(x) - \beta(0) + \beta(0) = \psi(x)$$

Thus  $u_g(x,t)+\beta(t)$  solves the original problem.

The bottom line now is that we have two problems to solve. The simpler problem

$$ho \, c \, eta'(t) = g(t)$$
 $eta(0) = rac{1}{l} \int_0^l \psi(x) \, dx$ 

can be solved immediately by integration:

$$\beta(t) = \int_0^t \frac{g(s)}{\rho c} \, ds + \frac{1}{l} \int_0^l \psi(x) \, dx$$

The main problem

$$\rho \ c \ \frac{\partial u_g}{\partial t} - \kappa \ \frac{\partial^2 u_g}{\partial x^2} = f(x,t) - g(t)$$
$$\frac{\partial u_g}{\partial x}(0,t) = \frac{\partial u_g}{\partial x}(l,t) = 0$$
$$u_g(x,0) = \psi(x) - \beta(0)$$

satisfies a compatibility condition, so we can expect to solve it by appropriate application of the Finite element method.